WALL CONFINEMENT EFFECT IN PARALLEL FLOW INDUCED VIBRATION

By SHAH NAWAZ MOHD. MALIK

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DEPARTMENT OF MECHANICAL ENGINEERING INDIAN INSTITUTE OF TECHNOLOGY KANPUR JULY 1974

WALL CONFINEMENT EFFECT IN PARALLEL FLOW INDUCED VIBRATION

A Thesis Submitted
In Partial Fulfilment of the Requirements
for the Degree of
MASTER OF TECHNOLOGY

By SHAH NAWAZ MOHD. MALIK

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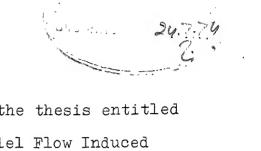
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POST GRADUATE OFFICE This thesis has been approved for the award of the transfer of Master of Table 19 and 19 regulations of the Inflant Institute of Technology and Dared. 6.8.74. 24

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CONTENTS

			Page
LIST OF FIGURES			(vii
LIST OF PLATES			(ix)
NOMENCLATURE			(x)
SYNOPSIS			(xi)
CHAPTER-I:		INTRODUCTION	
	1.1	:General	1
	1.2	:Previous Work	3
	1.3	:Present Work	18
CHAPTER-II :		EXPERIMENTAL SET-UP	
	2.1	:General	21
	2.2	:Components Fabricated	
		2.2.1 : Diffuser Assembly	23
		2.2.2 : Perspex External Flow-Tube	23
		2.2.3: Test-Rod Supports and Fixing Pins.	24
	2.3	:Test-Rod and Flow-Tube Dimension	25
	2.4	:Experimental Procedure	26
		2.4.1 : Test-Rod Preparation	28
		2.4.2 : Strain-gage Circuit	29

_		•
1	77	•
١.	v	

		2.4.3	:	Damping Measurement	30
		2.4.4	:	R.M.S. Response Measurement	31
		2.4.5	:	Measurement of Response Spectrum Components.	32 1
	2.5			Specifications of trument Used.	32
CHAPTER-III	•	RESUL	TS	AND DISCUSSIONS	
	3.1	:Gener	al		42
	3.2		aĪ.	Factor in the ly Filled External be.	42
		3.2.1	:	Graphical Plots	43
		3.2.2	:	Discussion	43
	3.3			Factor at Various Locities	
		3.3.1	:	Graphical Plots	45
		3.3.2	:	Discussion	45
	3.4			Response at Various Locities	3
		3.4.1	:	Graphical Plots	48
		3.4.2	:	Discussion	48
		3.4.3	:	Graphical Plots	50
		3.4.4	:	Discussion	50
	3.5	:Respon	se	Spectrum	
		3.5.1	:	Graphical Plots	51
		3.5.2	:	Discussion	51

CHAPTER-IV	:	CONCLUSIONS AND SCOPE FOR FUTURE WORK			
	4.1:	Conclusions			
		4.1.1: Effect of Water Level on the Damping	70		
		4.1.2: Effect of Flow Velocity on the Damping	y70		
		4.1.3: Effect on the R.M.S. Response	71		
		4.1.4 : Effect on the Power Spectral Density	72		
	4.2:	Scope for Future Work	72		
REFERENCES	:		74		
BIBLIOGRAPHY	•		78		
APPENDIX-1	:	Equation of Motion as Obtained by Chen and Wambsganss (23)	80		
APPENDIX-2	:	Part (A) : Diffuser Dimensions	84		
		Part (B): V-Notch Characteristics	84		
APPENDIX-3	:	Typical values of R.M.S. Response	85		

LIST OF FIGURES

	Page
Fig. 1 : Line Sketch of Test Set-Up.	38
Fig. 2 : Line Sketch of Flange Assembly for Flow-Tube.	39
Fig. 3 : Line Sketch of Test-Rod Suppor	ts. 40
Fig. 4 : Electrical Circuit Diagram for the Test Data Measurements.	41
Following Figures show the Wall confinement Damping Factor at various water level	nt effect on
Fig. 5 : For 3/8 inch Test-Rod.	54
Fig. 6 : For 1/2 inch Test-Rod.	55
Fig. 7 : For 5/8 inch Test-Rod.	56
Following Figures show the Wall effect on at various Flow Velocities	Damping Factor
Fig. 8 : For 3/8 inch Test-Rod.	57
Fig. 9 : For 1/2 inch Test-Rod.	58
Fig.10 : For 5/8 inch Test-Rod.	59
Following Figures show the Wall confinement the R.M.S. Response (in Log-Log scale).	nt effect on
Fig.11 : For 3/8 inch Test-Rod.	60
Fig. 12 : For 1/2 inch Test-Rod.	61
Fig.13 : For 5/8 inch Test-Rod.	62

Page

Following Figures show the Wall effect on the R.M.S.Response (in Linear, Scale) based on the Data of Line of Best Fit in Fig. 11, 12 and 13.

Fig.14 :	For 3/8 inch Test-Rod.	63
Fig.15 :	For 1/2 inch Test-Rod.	64
Fig.16 :	For 5/8 inch Test-Rod.	65
	Figures show the Power Spectral Density Flow Velocities.	plots
Fig.17 :	For 1/2 inch Test-Rod mounted in 2.0 inch Flow-Tube.	66
Fig.18 :	For 1/2 inch Test-Rod mounted in 2.5 inch Flow-Tube.	67
Fig.19 :	For 3/8 inch Test-Rod mounted in 2.0 inch Flow-Tube.	68
Fig.20 :	For 3/8 inch Test-Rod mounted in 1.0 inch Flow-Tube.	69

LIST OF PLATES

PLATE 1 : Test Set-Up (Front View).

PLATE 2 : Test Set-Up (Downstream-Side-View).

PLATE 3 : Instrumentation.

PLATE 4 : Perspex Test-Section.

NOMENCLATURE

Ø : Test Rod diameter (outer) in Inches.

t : Wall thickness of the Test Rod in Inches.

 Ψ : External Flow Tube diameter in Inches.

 \mathcal{A}_{W} : Damping factor with partially filled Flow-Tube.

JNW : Damping factor with no water in the Flow-Tube.

5 NF : Damping factor at no - flow and full water in the Flow-Tube.

X_D: Vertical_R.M.S. response (in Milli-Volts) of a test-rod in various Flow-Tubes.

: Vertical R.M.S. response for the largest diameter Flow - Tube (here 2.5" ψ) at a velocity at which X_D is measured.

d: Diameter Ratio \emptyset/ψ , a measure of wall confinement.

 ω : Damped Natural frequency of the Test Rod in ${\rm H}_{\rm Z}$, measured from the damping plots.

Stiffness: Ratio of stiffness or mass of the Test-Rod and mass ratio in question to the smallest size Test-Rod.

SYNOPSIS

of the
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Parallel Flow Induced Vibration

Submitted in Partial Fulfilment of the Requirements for the Degree

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SHAH NAWAZ MOHD. MALIK

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An experimental investigation has been made to obtain the effect of wall confinement on the vibration of elastic cylinder held in parallel fluid flow. Four kinds of vibration parameter are studied in an open loop water tunnel with variable circular flow area (in steps): effect of partial water-level on the damping factor, effect of flow velocity on the damping factor, effect of flow velocity on the amping factor, effect of flow velocity on the r.m.s. response and the effect on the power spectral density of response.

Experiments are carried out on aluminium tube with fixed-fixed end conditions. Four electrical resistance strain gages were fixed at half length to sense the signal in two perpendicular directions. Flow velocity was varied upto 60 feet/sec. Three different diameter test rods and external flow tubes were used.

Plots are given for different parameters as a function of flow velocity for each test rod in various external flow tubes. Comparison with available experimental results is made wherever possible.

The effect of wall confinement is to increase the various parameters studied here. But the rate of increase is more for the test rod with very low stiffness. However, if the two test rods have comparable stiffness the wall confinement has greater effect on the one with larger diameter (even if this test rod has greater stiffness).

CHAPTER-I

INTRODUCTION

1.1 General:

Analysis and prediction of dynamic behavior of fluid coupled elastic system has a very wide range of applications. We can classify them broadly in three categories:-

- (1) Fluid-elastic instability of cylinders in cross-flow.
- (2) Stability of an elastic tube conveying flowing fluid.
- (3) Fluid coupled elastic stability of cylinders in parallel flow.

All the three problems of dynamic interaction have very important roles to play in the Nuclear Reactor. Core components design and performance, for example, vibration of heat-exchanger tubes in cross-flow, bending vibration of piping systems and the response of fuel rods in parallel coolart flow to extract the heat generated during the controlled nuclear fission process. Both in the Fast Breeder as well as in the Thermal reactors, high

velocity of the coolant is the primary source of energy inducing the above mentioned instabilities.

Among the three, first one is the oldest problem of hydrodynamic instability predicted by Th. Von Kármán, while the second one was identified as early as in 1950 by Ashley and Haviland (1). But attention to the parallel flow induced vibration was given very recently during the investigation of reactor core components failure.

In the nuclear reactor core, series of uranium fuel pellets are contained in a metal cladding, forming a fuel rod. Cladding material is a must for several reasons, viz., retention of fission products, protection of the fuel from the coolant, facilitation of insertion and removal of fuel from the core and the exact location of fuel within the core. Several of these fuel rods forming a fuel cluster are fed into the coolant channel. The axially flowing coolant induces transverse vibration of fuel rods. Though the amplitude of vibrations are very small but the close spacing of fuel rods create the problem of wear, fatigue and fretting.

It has been found that vibrating fuel rods may also lead to incorrect flow distribution and the Neutron-flux oscillations, resulting in random perturbations in

the reactivity and kinetic behavior of nuclear system (2).

In the present work, the flow induced vibrations of elastic cylinder in parallel flow will only be dealt with.

1.2 Previous Work:

First study of the parallel flow induced vibration was carried out by Burgreen, Byrnes and Beneforado (3) while studying the problem of neat transfer to the water flowing parallel to a bundle of rods simulating a watercooled reactor core. Test were conducted on a vertical bundle of 13 and 37 rods in a closed flow loop. were collected for the mid point rod response and natural frequency at different flow velocities, wising strain gage circuitory. Measurements revealed that the vibration occur at all the flow velocities and the amplitude increases with the velocity. Frequency of vibrations in water were lower than in still air and remained almost constant for a wide range of water flow velocity. These facts led to the conclusion that the oscillations were of self excited type rather than forced one (for example, due to Karman vortices). Also, an analytical model was proposed using the force barance among flexural restoring force, damping and inertia force. Hydrodynamic force of water was identified as the forcing function. An emperical

relation for the response amplitude was given, using nondimensionalized parameters. Finally it was concluded that the parallel flowing fluid can excite oscillations at the natural frequency or more precisely at the damped natural frequency of the rods.

During the late '60s, number of flow induced vibration problems in reactor core components were observed notably in the reactors EBR II (4), GETR (5,1*), Big Rock-Point (6,2,3 and 4), Yankee (7, 5), SPERT III (8,1) and MSRE (9) etc.. These involved the problem of thermal power oscillations, fatigue, wear, mechanical failure of incore components and the neutron-flux oscillations. Though the occurance of such vibration problems were checked by redesign using trial and error but the mechanism of energy extraction and vibrational response propagation was not established clearly.

After Burgreen's work several experimental studies were conducted to predict the amplitude of vibration emparically in the United States (6 to 10), Sweden (11 to 13), France (14) and Canada (10, 11). In 1962, Quinn (7)

^{*} Digits with a bar represent the papers reported in bibliography.

Digits without bar represent the papers referred in the list of references.

performed tests on single and multiple rod assemblies both in one and two phase flow. A mathematical model was given and its solution was compared with the experimental data. However, the mathematical model contained several vaguely defined parameters and was difficult to apply.

Paidoussis (10) collected the experimental data of Quinn along with those of Swedish and French investigators and tried to correlate with the Burgreen et.al. emperical expression. He found deviations upto 10 times for water tests and upto 100 times for super heated stream tests. So, he himself derived the equation of motion for the transverse vibration of a beam in axial flow using the force balance technique involving the following forces:

- (1) flexural rigidity of the beam
- (2) longitudenal tension
- (3) force due to inertia
- (4) lateral fluid forces which are
 - (i) reaction on the cylinder of the force required to accelarate the fluid around it (12),
 - (ii) frictional normal drag force.

However, no account was given to the viscous and structural damping of the material.

Instead of solving the resulting equation of motion, Paidoussis gave an emperical relation for the mid rod displacement in terms of non-dimensionalized parameters of the equation of motion. Exponent of each parameter was found by keeping the other parameters fixed and plotting the vibration amplitude with respect to the parameter in question. He admitted that a better correlation can be obtained using the variables involving cross-flow, swirl velocity and surface roughness which were neglected owing to the absence of experimental data.

Further work of Paidoussis (11, 13, 14) deals with the solution of the governing equation of motion in the form of Y (ξ) e , where ω is the dimensionless complex frequency. The solution for the complex frequency ω , for the given system parameters, as a function of non-dimensionalized flow velocity was plotted on the Argand Diagram. It was predicted that system will be stable when the imaginary component of ω , $I_m(\omega) > 0$ and unstable when $I_m(\omega) < 0$, $I_m(\omega) = 0$ gave the condition of neutral stability.

The stability in the three lowest mode was considered. The stability diagrams indicated that the motion of cylinder will be damped out at lower flow

velocities (Real component of ω , $R_e(\omega)$ decreases) and at higher flow velocities system becomes unstable. However the first mode instability was of buckling type $(I_m(\omega) = R_e(\omega) = 0)$ and the higher mode instabilities were oscillatory $(R_e(\omega) \neq 0)$. These trends were almost similar for all types of boundary conditions viz. pinned, clamped and clamped-free.

established experimentally using cast rubber cylinders in a water flow-loop. Critical flow velocities for buckling and higher modes were measured and found to be close to the predicted analytical values. However, these hydro-elastic instabilities occured at very high flow velocities not to be encountered in the present day reactor technology. These "sub-critical vibration" were postulated to be due to the cross-flow component of axial flow velocity produced owing to non-uniform flow passages, secondary circulation, turbulence and separation behind the structures supporting the cylinder. It was anticipated that even in the smooth flow conditions (with White Noise spectrum) cylinder will extract energy from the noise spectrum corresponding to its normal mode and will be set into motion.

In his later work (15), Paidoussis modified the emperical relation predicting the mid rod displacement

on the basis of experiments on single cylinder with different mass, length, flexural rigidity, diameter, different flow conditions and method of supporting the cylinder. It was found that the swirling motion will increase the amplitude by a factor of 10 to 100 times above the predictions. Experiments with a bundle of 19 cylinders indicated the increase in vibration amplitude by a factor of 50 to 100, even in the quite flow conditions. However, at higher flow velocities this factor gradually reduced to a value of 5 to 10 with respect to the emperical predictions.

Analytical work carried out by Burgreen et.al (3)
Quinn (7) and Paidoussis (10) etc. identifies the viscosity of flow medium as the major factor affecting the
vibration, whereas the experimental results do not conforit precisely. Y.N. Chen (32) postulated the turbulence
in the flow as the major factor exciting the vibration.
He considers the vibration to be the self excited one
because the turbulence in the flow will cause rod oscillations which in turn will further increase the turbulence in the boundary layer. He formulated the dynamical
equation involving the longitudinal surface drag force,
inertia force of the rod, normal drag of the lateral flow
and the centrifugal force of the flow on the rod (due to

its curvature). Inertia force due to Coriolis accelaration was neglected on the assumption of smaller angular velocity having mixed derivative term $(\frac{\partial^2 y}{\partial x \partial t})$. Axial flow velocity was assumed to be the sum of mean flow velocity and a fluctuating component due to the turbulence. Governing equation of motion of the system was then simplified to the Mathieus type differential equation. Stability criteria was established on the basis of critical flow velocity that results in the buckling of the rod. As the turbulence in the flow has a broad frequency spectrum, so a degree of instability of the rod will always be excited. But the instability may be easily excited if the strength of turbulence is high at a frequency near the rod natural frequency.

Experimental data of Burgreen (3), Quinn (7), Sogreah (14), Pavlica and Marshall (10), Roström (11) were used to find the functional dependence of non-dimensionalized vibration amplitude on the flow velocity. Exponent found in this manner agreed very well with the Chen's theoretical predictions. Finally it was concluded that the rod vibration will not grow to a great extent because of the unsteadiness in the phase of flow velocity fluctuations.

During the late '60s, much progress in the flow induced vibration field was achieved. The concept of probabilistic and random nature of pressure field was first introduced by Reavis (16) and was used later by Gorman (17, 18), Chen and Wambsganss (21, 22, 23). Reavis analysed the first mode response of a pin-ended beam subjected to random pressure. Major assumptions were,

- (1) pressure field of turbulent air flow in a duct was assumed to be identical to that of water flowing over a cylinder
- (2) beam supports were assumed to be motionless and have no effect on the homogeneous pressure field. Comparison of theoretical values of the response with experimental results of Burgreen (3), Quinn (7), Pavlica and Marshall (8) and SOGREAH (14) indicated the experimental data to be higher by a factor ranging from 3.5 to 420, with 14 a representative value. He finally attempted to find the average disparity ratio between theoretical and experimental results and concluded that the lateral and longitudinal correlation density of random pressure field should be modified.

Using the results of Reavis (16) and wall pressure correlation data of Bakewell (19); Gorman (17) gave modified analytical expressions for the longitudinal and

peripherial spatial correlation densities on the basis of his own experimental measurements. Damping coefficient of the rod was determined at various flow velocities. Comparison of experimental and analytical predictions of r.m.s. displacement (in a vertical plane) indicated good agreement (maximum difference of 20% was found). This led to the justification that the vibrations are predominantly due to pressure fluctuations in the turbulent boundary layer around the rod. Motion of the rod end supports was also identified as a source of input energy. However, the leading draw-back of his measurements was the assumption of the same boundary layer properties at the surface of test element and at the external flow tube wall.

Further work of Gorman (18) deals with the vibration of cylinder in two phase flow using air and water mixture. In two phase flow, pressure fluctuations were attributed to the change in lateral momentum flux of the fluid as these are reflected from the fuel rod, unlike the one phase flow where it was due to the lateral movement of eddies. Measurements were taken for longitudinal and prepherial correlation coefficients, damping plots and r.m.s. rod response at various flow velocities and quality of air-water mixture. With the increase in simulated steam quality, displacement increased and a peak was obtained

at 12% steam quality followed by a gradual decrease in response amplitude. Longitudinal correlation coefficient increased rapidly with the steam quality to reach a first peak, it falls of f some what and then further increased with the steam quality. However, the peripherical correlation remained almost constant with the increase in steam quality. The damping coefficients were found to be about 4 times higher than those for one phase flow and exhibited the linear damping law, but the damping was almost independent of the steam quality.

Gorman's next attempt (20) was on a bundle of seven rods in two phase flow. Central rod was instrumented while the other six rods were used to simulate the actual flow and formations of sub-channels. Analytical model was further simplified, by using the result of his earlier work (18), to express the r.m.s. response explicitly as a function of power spectrum of the pressure drop corresponding to the rod fundamental frequency. Both the energy input as well as the response in a bundle was found higher than a single rod. Experimental data also revealed that for different mass flow rate, vibration amplitude attained a peak value at 16% quality and then either fell-off or held constant. Motion of the central rod in a bundle was purely

random and did not exhibit any preferred plane of oscillation. Though it was assumed that the random pressure field around the central rod in a bundle is homogeneous, but no justification was given for the basis of thinking so.

The first rigorous treatment of the problem was given by Chen and Wambsganss (21) using the Reavis (16) approach on the equation of motion given by Paidoussis (13). The statistics of boundary layer pressure field was based on Bakewell's (19) measurement on a body of revolution. Vibrations were postulated to be induced mainly due to near field turbulent boundary layer pressure fluctuation. Terms for damping (associated with the Internal viscoelastic effect and external viscous effect) and random pressure were included in the governing equation of motion* for the rod with simply supported end condition. presence of mixed derivative Coriolis force term, the system does not possess classical normal modes. equation of motion was simplified by neglecting the Coriolis force, damping and variable coefficient terms to obtain a set of orthogonal eigenfunction as the solution of resulting equation which was coupled in nature. Decoupling was established by neglecting the off diagonal terms resulting in a set of equations representing the

^{*} Please see Appendix-1.

normal modes of beam vibration. Power spectrum of the rod response was found with the help of system response function, spectral density of boundary layer pressure field and the joint acceptance (representing the effectiveness of the pressure in exciting a particular mode of the rod vibration).

Expressions for the statistics of turbulent wall pressure spectrum were found out, based on the analytical model of Corocos (24) and experimental data of Bakewell (19), Clinch (25) and Willmarth and Wooldridge (26). It is worthwhile to note that Willmarth and Wooldridge found the large energy density of the boundary layer pressure at low frequencies. Parameter study to find the influence of each parameter on the rod response was conducted for the range encountered in the present day reactor technology. It was suggested that the r.m.s. response is proportional to the flow velocity with a power of 1.5 to 2.0. Comparison of the exponents of various parameters with that of Burgreen et. al. (3), Paidousais(15) exhibited close similarities.

In the later work, Chen and Wambsganss (23) modified the earlier analytical model to accommodate any

type of boundary condition and gave expression for the effect of adjacent rod or proximity to the wall on the virtual added-mass of the rod. Using the force balance method they derived the equation of motion. However the exact solution was difficult to be found out owing to three reasons:

- (1) problem is non-self adjoint type
- (2) system does not possess classical normal modes
- (3) equation of motion contained variable coefficient term.

The resulting equations, after neglecting the Coriolis force term $(\frac{\delta^2 y}{\delta x \delta t})$, damping, variable coefficient and the terms associated with $\delta y/\delta x$, were used to find an adjoint system of equations such that these have same eigenvalues and the two set of eigenfunctions are biorthogonal. Modal expansion technique with Galerkin's method was used on the governing equation of motion in conjunction with the set of biorthogonal eigenfunction to obtain a set of ordinary differential equations. Decoupling of the system of equations was obtained by neglecting the off-diagonal terms on the basis of their smaller values. Solution for r.m.s. displacement was obtained in terms of system response function, power spectrum of random pressure and the longitudinal cross correlation of random pressure.

Expressions for power spectral density of pressure field were modified on the basis of new experimental data of Wambsganss and Zaleski (22). Damping in a particular mode is the sum of three terms. viz.

- (1) damping associated with internal viscoelastic and external viscous effects
- (2) damping contribution of normal drag force
- (3) damping due to Coriolis force.

The later two damping effects increase with the flow velocity. However, the damping contribution of Coriolis force depends on the end conditions. If the rod ends are not allowed to move (eigen value problem becomes self-adjoint), the Coriolis force ceases to give any contribution to damp the motion. If the rods are elastically supported or free at one end, the Coriolis force damping contribution becomes maximum surpassing the other damping effects.

Experimental investigations were carried out to find the dependence of fundamental frequency and damping coefficient on the flow velocity by applying external harmonic excitation using an electromagnet assembly. Comparison of experimental data with the analytical results indicated that for fixed-ends the frequency decreases with

the increase in flow velocity and the rod may be subjected to buckling type instability at the critical velocity while for the cantilevered rods frequency increases with the flow velocity and rod may be subject to flutter type instability; both the damping and fundamental frequency plots displayed a close agreement with the analytical results.

Plots of r.m.s. rod displacement at the mid
point of fixed-fixed rod and at the free end of cantilevered rod displayed a wide gap with respect to analytical solution at low flow velocities (maximum difference
was as high as 300%). The lower values of prediction
were attributed to far field acoustic noise and structural
borne excitation which were not included in the analytical
model. Hydraulic diameter was also identified as an active
parameter. Lower hydraulic diameter means lower turbulence level resulting in a smaller r.m.s. response. The
dependence of r.m.s. response of flow velocity was approximated by a power function relationship. The exponent
depends on damping, hydraulic diameter, flow noise and
rod support conditions.

However, no experimental attempt was made to find the effect of adjacent rod or proximity to the wall on

the response, damping and natural frequency behavior.

Hence the effect of virtual mass on the response is yet to be explored.

The latest work of Paidoussis (34, 16) also takes into account the effect of gravity and external pressure on the response of cylinder. These two effects are prominent if the cylinder is held in vertical configuration and the system mean pressure is quite high. Argand diagrams were modified to represent the stability of the cylinder, when the above two effects are included.

1.3 Present Work:

In the available literature no work (experimental, emperical or analytical) has been reported on the effect of coolant flowtube wall on the response of vibrating cylinders. The present work has been initiated with an aim to understand experimentally the adjacent wall effect on the flow induced vibrations of flexible cylindrical rod in single, phase (liquid) axial flow.

Like the previous investigators, here also it was assumed that the vibrations of elastic cylinder are excited mainly due to random pressure fluctuations in

the turbulent boundary layer around the lateral surface of cylinder. Other possible sources of excitation are the motion of end supports and the pump flow pulsations etc.

Open loop water test-rig was designed, fabricated and installed to achieve water velocities well
above the values encountered in the present day reactor
technology. Test set-up has also the facility for
varying the cross-sectional area of external flow tube
(in steps) into which different diameter elastic cylinders are to be tested.

Strain gage circuitory has been used to find the mid rod response after proper amplification of the weak signals. Damping plots were recorded on the oscillograph both in full flow at different water velocities and partially filled flow tube with various water level (for example 1/2 filled, 3/4, 7/8 filled etc.). Owing to unknown distribution of forcing function (random pressure) it was difficult to convert the strain gage out-put into the displacement statistics; so, the strain values were used directly to represent the response of test rod. The rod end support conditions are fixed-fixed. Test data were also collected for the r.m.s. values of response signal using the random signal r.m.s.

voltmeter. Frequency response characteristics of the associated power spectrum was measured with the help of a band-pass filter by, varying the central frequency of filtered input to the true 'r.m.s. voltmeter.

Finally the curves were plotted for the above mentioned parameters of damping factor, r.m.s. values of response and power-spectrum as a function of mean axial flow velocity for different diameter of external flow tube as well as the test rods. The effect of adjacent boundary wall of the flow tube on the response of the test rods was discussed from these graphs.

CHAPTER-II

EXPERIMENTAL SET-UP

2.1 General:

In this chapter details have been given for the test set-up, components fabricated and experimental pro-The schematic diagram of the experimental set-up is presented in figure (1). Test set-up was designed to provide open-loop water circuit with variable flow velocity. Centrifugal pump used to supply the water has a capacity of 1000 gallons/minute at a head of 60 feet with 6 inch discharge pipe diameter. Since at high velocities major portion of head loss is due to friction losses in the pipe and the loss in 90 degree bends and elbows. care has been taken to provide these in the low velocity regions, that is, in 6 inch diameter cross-sections. Because of lack of space in the pump room, the test section was installed on a separate flume and was connected to the pump by a 6 inch pipe line. To mak the flow uniform, honey comb was fixed in the flow area at the upstream side, Diffusers with semi cone angle of 5 to 10 degrees were

used to reduce gradually the cross-sectional area from 6 inch to the desired test-section diameter*. further the flow to be uniform and fully developed, G.I. pipe of length greater than 40 diameters was connected to the upstream diffuser. This pipe length is connected to the perspex test section through flanges and bolts with rubber packing. Choice of perspex for the external flow tube was made owing to its high flexibility and ease in simulating actual conditions even at very low velocities, Aluminium tubes were selected as the test-rods due to their lower stiffness resulting in larger deflections. Test rod was supported axially into the external perspex flow tube through four mild steel pins at each end. At the downstream side of the test section a diffuser is again connected to increase the cross-section gradually to 6 inch diameter. Finally a 90 degree elbow is used to discharge the water vertically in the flume. Discharge from the flume was found by measuring the head over the V-notch. interval of about 5 minutes was needed to establish constant head over the V-notch for a particular gate valve opening. Mean-axial-flow velocity of water in the test section was calculated from the discharge values. Enough care was taken to make all the joints leak-proof.

^{*} Please see the Appendix-2

2.2 Components Fabricated:

In the present work parts fabricated involved mainly the following components:-

diffuser assembly, perspex external flow-tube assembly, test rod supports and fixing pins, flanges for 6 inch gate valve and pipe threading on various diameter G.I. pipe line.

2.2.1 Diffuser Assembly:

of the external flow tube, diffusers were made out of 1/8 inch thick mild steel plate and were hot rolled manually to conical shape. One end was welded to 6 inch diameter, 1 foot length G.I. pipe and the other end to 6 inch length of the desired flow-tube diameter G.I. pipe. Finally the flanges were mounted at the two threaded end of diffuser assemblies. To take out the strain gage lead wires, a brass nipple was brazed to the smaller diameter side of down-stream diffuser.

2.2.2 Perspex External Flow+Tube Assembly:

Flanges were made, using perspex sheet of 1 inch thickness. Flanges were bonded to the perspex tube with the help of acetone. Hub of the flange was also used to

fix the end supports of test rod using 4 mild steel fixing pins on its lateral surface. Details of a typical flange assembly for 2 inch diameter flow tube is given in the figure (2).

2.2.3 Test Rod Supports and Fixing Pins:

Aluminium test rod was supported at both the ends by means of test rod supports. These in turn were supported to the external perspex flow tube with the help of four 1/8 inch diameter mild steel fixing pins. of the test rod supports, made of mild steel, are given in the figure (3). It has three main parts: 1 inch length of diameter d_1 which press-fits in the internal diameter of aluminium test rod; section d, was joined to 1/2 inch diameter uniform section by the tapered portion; axial length of the uniform portion was kept 7/16 inch. Two 1/8 inch tapped holes at 180 degrees to each other were made along with a through hole of 1/8 inch diameter having 3/16 inch axial separation with respect to the tapped holes. Close dimensional accuracy was maintained to match these holes with those on the perspex flow tube mentioned in section 2.2.2. Finally the upstream and downstream end was made stream lined to avoid flow separation. In the down stream rod support, central hole was drilled

and connected to a drilled hole on its lateral surface to take out the strain gage lead wires coming through the annulus of the test rod.

To support the test rod structure in the external flow tube, two 1/8 inch fully threaded pins were screwed in the tapped holes of the rod support and a long mild steel pin threaded at one end was rassed into the through hole and screwed on the outer surface of perspex flow tube. In this way, fixing at 90 degree to each other in both the horizontal and vertical plane was achieved. The above mentioned precision work was carried out with close tolerances of the order of 0.005 inch.

2.3 Test Rod and Flow Tube Dimension:

In the present study of the adjacent flow tube wall effect on the response of test rods, following dimensions of the two were selected:

- (1) Perspex flow tubes of internal diameter 2.5 inch, 2.0 and 1.0 inch with 1/8 inch wall thickness and 4 feet length.
- (2) Aluminium hallow test rods of outer diameter 5/8 inch, 1/2 and 3/8 inch with the wall thickness of 0.049 inch, 0.0595 and 0.039 inch respectively and 46 inch length.

2.4 Experimental Procedure:

As outlined earlier, an elastic cylinder in axial flow will be set into oscillations, due to the random pressure distribution, in the turbulent boundary layer, upon its lateral surface. Pressure distribution around the oscillating cylinder will be unsteady and non-stationary in nature because of departure from purely axial; uniform and steady flow. The random forcing function on the test rod will result in the ramdomness of the response. Though the earlier investigators had identified the vibration to be of random in nature but they had attempted analytical solution as well as experimental measurements of response amplitude only in one particular plane, for example, vertical component. Whereas in actual condition the amplitude of vibration will not be having any preferred plane of oscillation. The actual response of the test rod can be found by measuring its component in two perpendicular planes.

There are several methods reported in earlier work for measuring the response of the test rod. Paidoussis (15) used sub-miniature piezoelectric accelerometer screwed to the mid span of the rod and integrating the output twice to obtain the displacement signal. Chen (23)

has used the optical displacement tracking technique. Electrical resistance strain gages were used by Burgreen et. al. (3) and Gorman (17). One other possible method, now in use, is based on the principle of electro-magnetic induction. Constant electro-magnetic field is applied normal to the lateral surface of the test rod and a coil is wound around the lateral surface of the rod where the response is to be measured, so an e.m.f. will be induced in a direction mutually perpendicular to the plane of oscillations and plane of magnetic field. This e.m.f. generated will depend on the number of turns in the coil and the rate of magnetic flux linking the coil. The voltage generated will represent the velocity of the test rod in a direction perpendicular to the rod axes as well as the magnetic field. However the system becomes very expensive because a very large magnetic field is needed to obtain measurable signal in a gap of 2 inch or more between the magnetic pole pieces to accommodate the external flow tube in between it.

In the present work the resistance strain-gages have been used because of their low cost, easier availability and good accuracy. To obtain actual response, strain was measured both in horizontal and vertical plane.

To minimise the curvature effect, due to small diameter of the test rod, gages of minimum possible dimension were selected.

2.4.1. Test-Rod Preparation:

Bending strain at mid span was sensed using four gages 90 degrees to each other. Gages were fixed on the lateral surface of the test rod, parallel to the axes.

Lead wires were then passed into the annulus of the test rod through the holes (1.5 m.m. diameter) on its lateral surface. These holes were drilled on the down stream side sufficiently far from the mid span to have negligible effect on the response. Lead wires were finally taken out of the downstream rod end support.

Water proofing of the strain gages is the crucial part of the test rod preparation. Gages have to with-stand following severe environmental conditions:-

- (1) Leakage of water, humidity, moisture and thereby absorbing water into the adhesive cement even at no water flow conditions.
- (2) Tangential viscous drag force on the test rod which is a function of flow velocity.

A method of double protection, after some initial gage failure, was finally used for water-proofing the gages.

A coat of melted cerese wax was applied over the gage and the adjoining area upto lead wire and gage soldered joint. Smoothness of the aluminium test rod surface in the vicinity of gage installation was reduced using emery paper and needle files to obtain proper bonding. liquid water proofing compound was then spread over the entire area including the cerese waxed surface using hair brush and acetone as thinner. Coating of the water proofing compound was done twice with an interval of 24 hours. Water proofing compound available in the domestic market had a tendency to absorb atmospheric air in the form of air bubles during its wet state. So, to avoid the possibility of holes and vents, double coat of water, proofing synthetic enamel paint was applied with 12 hours intermediate drying time. Thickness of the total gage preparation was less than 1/8 inch.

2.4.2 Strain-Gage Circuit:

Both the strain gage pairs in horizontal and vertical direction were connected separately to two bridge balancing unit. Gages were connected to two adjacent arms of each balancing unit in half bridge mode. This will result in cancellation of signal due to temperature and far field (acoustic) noise of the system. As outlined

by Wambsgnss (22), far field noise in the external flow, tube concentric with the test rod will give rise to the equal pressure contribution around the circumference of the test rod at a particular cross-section and at a particular instant. Hence, the response signal will be only due to near field flow noise of turbulent eddies. To amplify the output signal, both bridge balance units employed had a built-in amplifier of gain 50.

2.4.3. Damping Measurements:

Total system damping will be the contribution of several factors viz., structural damping of the test rod material, external viscous damping, damping due to normal drag force and Coriolis force. Damping due to these factors increases with the axial flow velocity except the structural damping. It is difficult to find the damping due to each factors separately. So, an attempt was made to find the equivalent viscous damping of the system. Initial deflection in horizontal gaged plane was given by attaching a weight to the steel wire (0.010 inch diameter) connected to the test rod. Diameter of the hole drilled in the perspex flow tube, through which the steel wire passes, was 0.012 inch to prevent any possible leakage. Damping plots at various flow velocities were

recorded on Visicorder Oscillograph by cutting the thread attached to the weight. Sufficient care was taken to remove the noise components in the oscillograph, which is quite susceptible, due to the strong magnetic field of galvanometer assembly.

To find the effect of various water level on the damping behavior, plots were also taken for partially filled external flow tube. Damping factor was calculated by measuring Logarithmic decreement in several successive cycles of damping trace.

2.4.4 R.M.S. Response Measurement:

R.M.S. value of the response, at various flow velocities, was measured by means of true r.m.s. voltmeter. Schemetic circuit is given in Figure (4). Output of the vertical and horizontal strain balancing units was fed to two D.C. amplifiers of variable gain. Output of each amplifier was then fed to the input of true r.m.s. meter. Gain of the D.C. amplifiers was adjusted to obtain measurable reading. R.M.S. value of both the signals was read separately each at a time. Because the random forcing function is unsteady in nature, so the r.m.s. meter pointer deflection was observed for 3 to 4 minutes and an average value was noted down. Proper grounding of the equipments was done to minimise the noise in the electrical circuit.

2.4.5. Measurement of Response Spectrum Components:

Power spetral density of the response at different flow velocities was measured to find the major contribution at various frequencies. For this purpose a constant band pass, variable frequency, filter was used in series with the true r.m.s. voltmeter. Both the horizontal and vertical strain signals were fed, one at a time, to the filter and the centre frequency was varied continuously. Major r.m.s. components were read on the r.m.s. meter, after proper amplification of total signal to the variable filter. Here also an average reading of the meter was noted down to cope with the unsteady nature of the response signal.

Measurements outlined in section 2.4.3, 2.4.4 and 2.4.5 were taken for all the test rods using various Size external perspex flow tubes.

2.5 Leading Specifications of the Instrument Used:

(A) <u>Visicorder Oscillograph</u>

Model : 906T, Honeywell Inc. U.S.A. Ultraviolet light recording

type.

Number of Channels: Six (active)

Galvanometers : M200-120, moving coil

miniaturé type.

Frequency : 0 to 120 Hz flat (± 5%) Response

Current

Sensitivity

: 25.5 Micro-Ampere/inch.

Voltage

Sensitivity

1.58 Milli-Volts/inch.

Deflection

Traces may overlap, permitting

a full 6 inch peak to peak

amplitude.

Record Speeds

Four speeds selectable ranges 5, 25, 50 and 100 millimeters

per second.

Recording Paper

Kodak Linagraph Direct Print

type (using Latent Image Intensification process).

(B) <u>Strain Indicator</u>

Model

Vishay/Ellis-10, Vishay Instruments Inc. U.S.A.

Function

Measurement of strain with half-bridge or quarter bridge strain gages.

Range of Resistance

: 100 to 1000 ohms.

Resolution

: + 2 Micro-Strain.

Range

: <u>+</u> 5000 Micro-Strain.

Bridge Excitation:

4.5 Volt D.C.

SCOPE Output

(i) External load 100 Kohms or more:-

 \pm 0.5 Volt swing.

± 5000 Micro-Strain, linear within 2%.

(ii) External load 200 ohms:-

+ 250 Micro-Ampere swing.

± 1000 Micro-Strain, linear within 5%. (iii) Amplifier Characteristics:-

Single-ended output.

Frequency: Flat between Response 0 to 500 Hz.

Output Impedence : 500 ohms for a dynamic swing

of + 1000 Micro-

Strain.

(C) <u>Variable Filter</u>

Model : 1-159-0001, Consolidated Electrodynamics Corpn. U.S. A.

Function : Constant percentage Band-Pass filtering with variable

center frequency.

Bandwidth : 3% (constant) between 3 dB points. Minimum attenuation one Octave either side of the tuned frequency is 50 dB

minimum.

Frequency Range : 8 to 2500 Hz (center frequency

tuning range).

Input Impedence : 50 Kohms + 2%

Noise Output : 1 Milli-Volts r.m.s. or less,

with input shorted.

(D) R.M.S. Meter

Model : 55A06 Disa, Denmark.

Function : Root-mean-square voltage

measurement of electrical signals irrespective of

wave form.

Frequency : 3 Hz to 200,000 Hz. Response

Crest Factor : 5:1

Input Impedence 1 Mega-ohm in parallel with

50 p.F.

Meter Response Low or high damped.

Input Voltage : 0.2 Milli-Volt to 100 Volts.

Range

Noise Level Less than 50 Micro-Volt r.m.s.

Common Mode 100 at 1000 Hz.

Rejection

(E) Solid State Differential Amplifier

TDA-875, Astrodata California Model

(U.S.A.).

Function Astrodata TDA series consists

of amplifiers that are designed for use in missile checkout systems, in low-level instrumentation systems and in online control systems. They are for use in all cases in which high Common Mode Rejection, extraordinarily high input

Impedence and high reliability

are essential.

Voltage Gain : 1 to 2000 (11 fixed steps).

Mode of Operation: Single ended as well as

Differential.

Input Impedence Greater than 100 Mega-ohms in

parallel with 0.0007 Micro-farad

Noise Level Less than 5 Micro-Volts r.m.s.

over the full frequency band.

Single Ended Mode Operation :-

Grain Stability DC to 1 Kc(within 0.1%).

Output Voltage + 10 Volts (maximum) from DC

to 50 Kc.

Load Resistance Minimum 200 ohms. :

Less than 1.0 ohm in series Output Impedence :

with 50 Micro-henries.

Frequency Response: + 1.0%, DC to 10 Kc.

Down 3 dB at 150 Kc.

Harmonic Less than 0.1% r.m.s., upto

Distortion 10 Kc.

Differential Mode Operation :-

+ 1.0% from DC to 300 Hz. Frequency Response

Down 3.0 dB at 5 Kc with

1000 ohms load.

Output Resistance: Less than 25 ohms

Load Resistance Minimum 1000 ohms.

+ 10 Volts (maximum) from Output Voltage

DC to 5 Kc.

Common-Mode Greater than 120 dB at 60 Hz

for 1000 ohms of line Rejection

unbalance.

Common-Mode 300 Volts - step common mode

does not cause overload: Level

DC to 1 Kc (within + 0.1% DC Gain

of the full scale) -Stability

Strain Gages (F)

> NKF-2, Rohits & Co., Roorkee, Model

India.

120 ohms (+ 1.0%) Gage Resistance

2.68 (+1.0%)Gage Factor

Strain Sensing Material

A Chromium-Copper alloy having traces of Si + Mn and it is recommended for dynamic strain measurements, because of its high temperature sensitivity and high gage factor.

Gage Dimensions

Grid length 2.0 millimeter Grid width 0.5 millimeter Base length 8.0 millimeter Base width 5.0 millimeter Hellical Grid.

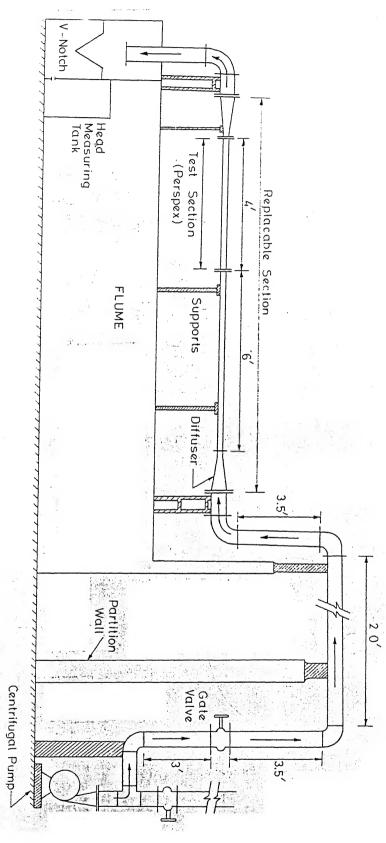
Maximum Current

40 Milli-Ampere.

Maximum Temperature Permissible

80 degrees centigrade.

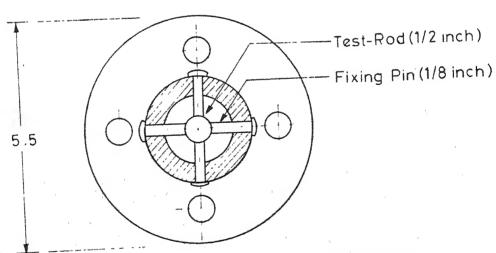
Cross-Sensitivity: 1.5%



Test Set-up

FIG .-1

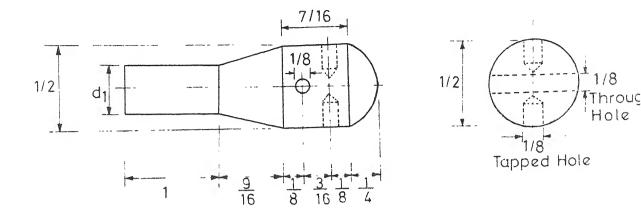
Loaded With Test-Rod



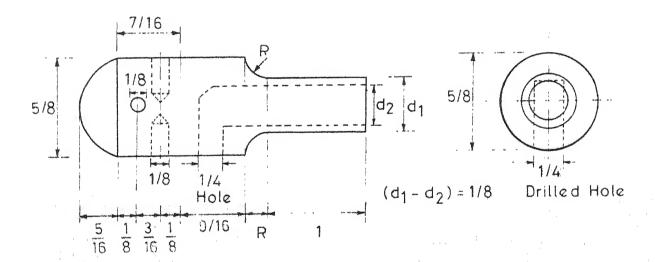
Material: Perspex All Dimensions In Inches

Flange Assembly for 2inch Flow-tube

Upstream Support



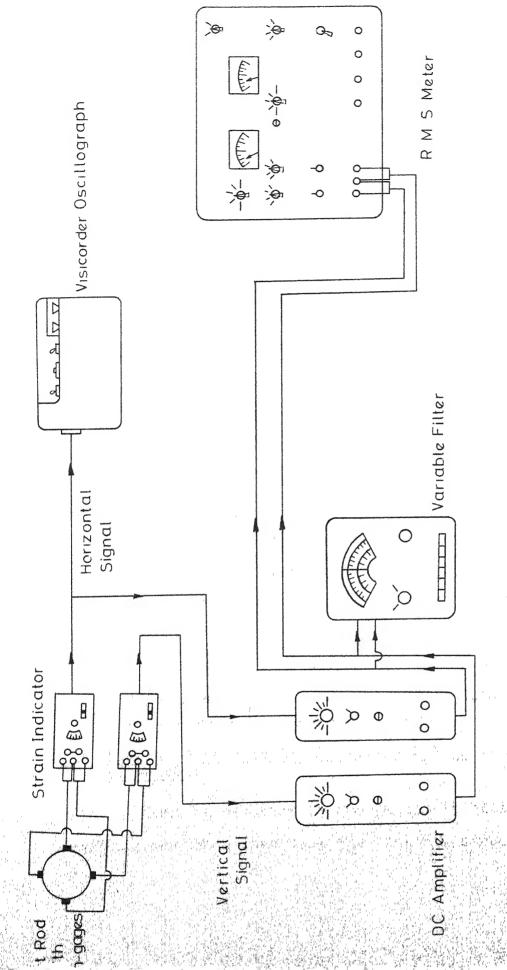
Downstream Support



Material: Mild Steel All Dimensions In Inches

Test-Rod Supports

FIG.-3



Electrical Circuit Diagram

F16.-4

CHAPTER-III

RESULTS AND DISCUSSIONS

3.1 General:

In this chapter, graphical plots and results have been presented for the wall effect on the following parameters:

- (1) Damping factor in the partially filled external flow tube.
- (2) Damping factor at various flow velocities.
- (3) R.M.S. response at various flow velocities.
- (4) Response power spectrum.

Theoretical stiffness and mass ratio for the three test rods are tabulated below:

Size	Stiffness Ratio	Mass Ratio
3/8 inch.	1	1
1/2 inch	3. 95	2.085
5/8 inch	6.6	2.18

3.2 Damping Factor in the Partially Filled External Flow-Tube:

Damping behavior in the partially filled tube is a complex phenomenon involving the flexural structure subjected

to two fluids (having different density and viscosity) in a confined elastic conduit.

3.2.1 Graphical Plots:

Figs. 5,6,7 show the wall effect on the damping of test-rod at various water level in the flow tube. Increase in the damping factor, represented as $\Im W/\Im W$, is taken as ordinate and the water level (0 to Fully-filled) as abscissa. Each figure was plotted for a particular size of test rod.

3.2.2 Discussion:

Damping values for no water, JNW, for each rod with fixed end conditions mounted in the different external flow—tube are fairly high (about 8 to 10 times the values of structural damping in a simple laboratory test.). This may be attributed to the end fixity conditions. Damping values SNW for a test rod in different flow tubes were also not same. This can be explained on the basis of the difference in rigidity of the rod supporting structure (which also include the perspex flow tube).

Fig. 5 shows the wall effect for the 3/8 inch diameter test rod. Damping increases with the water level. Even for 2.5 inch flow tube, there is a 55% increase in damping from zero to full water level. For 2.0 inch flow tube, water level above 3/4 has a faster rise in damping.

1.0 inch flow tube (area of cross-section being 1/6.25 of the 2.5 inch flow tube) gives the steepest rise in damping values upto 215%. We can infer that the wall effect is small for larger flow tubes filled upto 3/4th level and above this level it is considerable. For the smallest flow tube (wall confinement parameter d = 0.375) the damping factor is very sensitive when the tube is partially filled.

Fig. 6 shows the same trend, for a 1/2 inch test rod except that the rate of increase in damping is very low. For example, the increase in damping is 20% for 2.5 inch flow tube and 35% for 1.0 inch flow tube at full water level. Percentage increase in damping for the 1/2 inch test rod is quite less compared with the 3/8 inch test rod. This is due to larger stiffness of 1/2 inch test rod.

Fig. 7 shows the wall effect for the 5/8 inch test rod. Here the trends are slightly different. Maximum damping values are attained at approximately 1/2 water level and after that the damping remains constant upto full water level. However, the total percentage increase in damping is almost same as to that of 1/2 inch test rod (that is 25% for 2.5 inch flow tube and 38% for 1.0 inch flow tube). The change in behavior can be attributed to the damping due to normal drag force (which is a function of test rod surface area and, hence, dominant here) attaining its maximum value when the water level reaches 1/2.

Test rod damped natural frequencies are also tabulated in Figs. 5, 6 and 7. Natural frequencies are almost constant for each test rod and are quite independent of the water level in the flow tube. This shows that the confinement has very small effect on the natural frequency of vibration.

A comparison of all the three graphs show that the damping will increase at a faster rate and the wall effect will be prominent if the initial damping factor with no water is small.

3.3 Damping Factor at Various Flow Velocities

3.3.1 Graphical Plots:

Figs. 8,9 and 10 display the wall effect on the damping factor of the test rod at various flow velocities. The percentage increase in damping factor, \$F/\$NF, is plotted as ordinate for the same test rod, when mounted in the different flow tubes, with mean axial flow velocity as abscissa. Damping plots were taken for flow velocities upto 20 to 25 feet/sec; above which the damping plots became quite random in nature.

3.3.2 Discussion:

All the three figures show a linear increase in damping factor with flow velocity. This fact is identical

with the experimental results of Chen and Wambsganss(23) who validated it theoretically also.

The damped natural frequencies of oscillation, measured from the damping plots, remained almost constant at various flow velocities as well as with wall confinement. These have been tabulated in Figs. 5, 6 and 7. The effect of flow velocity on natural frequency is identical in nature with the previous results of Chen(23) and Burgreen et al(3). Chen found a decrease in frequency of order of 4 to 5 % for a velocity range of 0 to 100 ft/sec. From the values given in Figs. 5,6 and 7 it can be inferred that the wall confinement has negligible effect on the natural frequency at moderate velocity range.

Fig. 8 shows the effect of wall on the damping factor for a 3/8 inch test rod. Wall effect is quite prominent even when the change is made from 2.5 inch to 2.0 inch externs al flow tube. For the 1.0 inch flow tube the percentage increase in damping is large (for example 40% for a velocity of 17 ft/sec). This trend is expected because the stiffness of this test rod is very small and the normal drag force increases with the flow velocity.

Fig. 9 gives the trend for 1/2 inch test rod.

Increase in damping with flow velocity and wall confinement

effect is small when the flow tube diameter is reduced from 2.5 inch to 2.0 inch. But, for the 1.0 inch flow tube the wall effect is appreciable. The overall increase in damping is about 16% to 24% as we go from 2.5 inch to 1.0 inch flow tube at a velocity of 25 ft/sec as compared to no-flow conditions. This relatively lower increase in damping with the flow velocity may be due to larger initial damping values (3NF) for 1/2 inch test rod as compared to other test rods.

Fig. 10 indicates the effect on a 5/8 inch test rod. Wall effect is small as we move from 2.5 inch to 2.0 inch flow tube; but it becomes most prominent for the 1.0 inch flow tube (55% increase in damping for a velocity of 17 ft/sec as compared to no-flow condition. This can be attributed to the fact that the initial damping values (3NF) are moderate and the stiffness is of the same order as that of 1/2 inch test rod, thus the motion has more chances to be damped out with the flow velocity. Also, the damping contribution due to normal drag force for a 5/8 inch test rod will be maximum.

A general inference can be drawn that if the initial damping is small then the wall confinement will have more effect on a lower stiffness test rod to damp out the motion as the flow velocity increases.

3.4 R.M.S. Response at Various Flow Velocities:

3.4.1 Graphical Plots:

Figs. 11,12 and 13 give the wall confinement effect on the r.m.s. response of test rods as a function of mean-axial flow velocity. Plots are given on the log-log paper with r.m.s. response in Milli-volts* as ordinate and flow velocity (ft/sec) as abscissa. Here, only vertical component of the response was plotted because in almost every case it was found that the horizontal component was around 90% of the vertical component** and the horizontal component varied in the same manner with respect to flow velocity. Thus the vertical component can represent the wall effect.

3.4.2 Discussion:

Experimental data fit very well in the power function relationship of the type $Y = A \cdot (X)^B$, with minimum standard

^{*} Strain-Indicator Scope-Output specification is 10 Micro-Strain per Milli-volt. The gain of D.C. Amplifier is selected as 100. Thus the output is 10 Micro-Strain per 100 Milli-volts.

^{**} Typical values of horizontal and vertical components of the response are given in Appendix-3.

error. Values of exponent B were computed and are tabulated below:

Exponent B

Flow tube size	1.0 inch	2.0 inch	2.5 inch
Test rod size			
3/8 inch	2.48985	1.76159	1.52096
1/2 inch	1.52461	1.27949	1.15267
5/8 i nch	1.87351	1.52034	1.36915

Power function relationship satisfied by the experimental data agree with the theoretical and experimental results of Chen and Wambsganss(21, 23). Chen(21) found the value of exponent B in the range of 1.5 to 2.0. We can infer, from the present experiment, that the wall confinement can increase the exponent B in the range of 1.15 to 2.5. However, the upper and lower bound for the exponent B will depend upon the flow noise in the test-loop, damping, rod end support conditions and flexural rigidity.

Wall effect for a 3/8 inch test rod is presented in Fig. 11. R.M.S. response increases considerably even for the confinement from 2.5 inch to 2.0 inch flow tube. Confinement effect for 1.2 inch size influences the response

level to a maximum value. This is due to the fact that the stiffness of the test rod is small and a disturbance created around the test rod will become more severe after reflecting from the flow tube wall, resulting in further increase in response.

Figs. 12 and 13 represent the effect on 1/2 inch and 5/8 inch test rod. The response level and its rate of increase is more for 5/8 inch test rod than that for 1/2 inch size, though the stiffness is large (theoretically 65%) for 5/8 inch test rod. This is attributed to the wall confinement effect. The confinement parameter d is 25% more for 5/8 inch test rod compared to 1/2 inch rod.

3.4.3 Graphical Plots:

Figs. 14, 15 and 16 represent the wall effect on the r.m.s. response (vertical) as a function of flow velocity in the linear scale. Straight lines of best fit in Figs. 11, 12 and 13 were used to represent the numerator of the ordinate denoted as X_D . Denominator (X_{LD}) of the ordinate is the vertical response values for the largest diameter flow tube.

3.4.4 Discussion:

Each figure show the wall confinement effect for various test rods on a linear scale. These figures express the conclusions of Section 3.4.2 in a much more lucid way.

The response and its rate of increase with the flow velocity decreases with the increase in stiffness of the test rods (compare Figs. 14 and 15 or 14 and 16). But at the same time, wall confinement effect helps in building up the response and the rate of increase in response for the larger stiffness test rod (Figs. 15 and 16). In the present work, larger stiffness test rods have also the larger outer diameter (that is the greater wall confinement). It will be interesting to work on the lower stiffness with larger diameter and constant stiffness with different diameter test rods to understand the wall confinement effect in a broader sense.

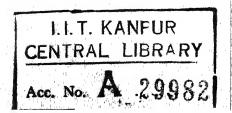
3.5 Response Spectrum:

3.5.1 Graphical Plots:

Power spectral density of the test rod response is plotted for different flow velocities in Figs. 17,18,19 and 20. Here also the vertical component of the response is selected as ordinate. Frequency in Hertz is taken as abscissa in a linear scale. To measure the spectral components accurately the gain of D.C. amplifier was further increased by 5 times.

3.5.2 Discussion:

A glance on these plots will reveal that, apart from pump (24.2 Hz) and the line frequency noise (50 Hz), a



considerable amount of energy is distributed at frequencies other than the rod natural frequency. The energy distribution is mainly at the rod natural frequency for lower values of flow velocities. But at higher flow velocities the energy distribution is less at the rod natural frequency. This behavior can be due to two reasons:

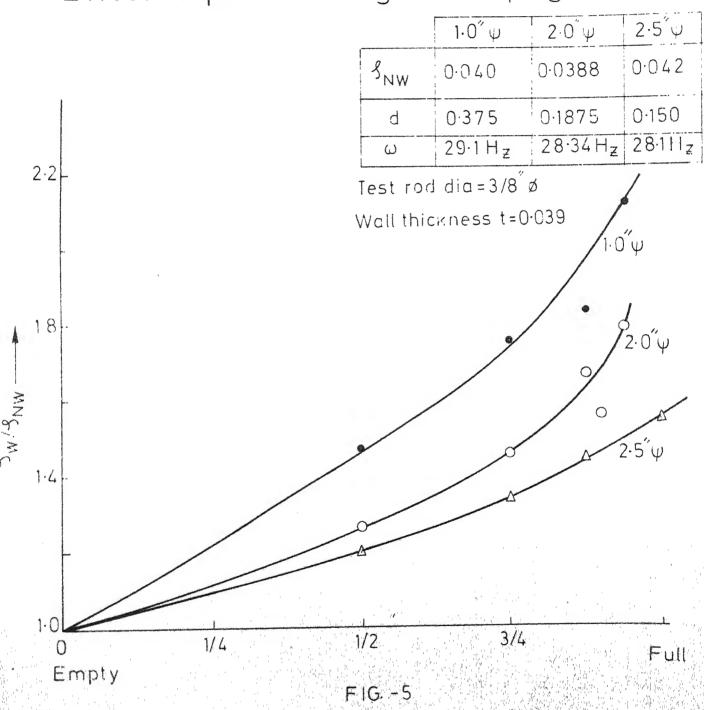
- 1) Test rods possess high damping factor, thus the frequency-response function of the test rod will be fairly broad and will contain additional frequencies, on the both sides of fundamental frequency, with significant response.
- 2) The spectral energy distribution of the pressure fluctuation in the turbulent boundary layer around the test rod increases at the lower frequencies with the increase in flow velocity. This is attributed to the fact that the energy of larger eddies will be transferred to the smaller ones, due to the turbulence decay process(32).

Figs. 17 and 18 show the wall confinement effect on the power spectral distribution for the 1/2 inch test rod.

A 25% increase in wall confinement (when the flow tube diameter is changed from 2.5 inch to 2.0 inch) increases the energy distribution at lower frequencies in a larger proportion than

its increase at the rod natural frequency. Figs. 19 and 20 display the same effect for a 3/8 inch test rod for a 50% increase in wall confinement parameter d (when the flow tube is changed from 2.0 inch to 1.0 inch size).

Effect of partial filling on damping factor



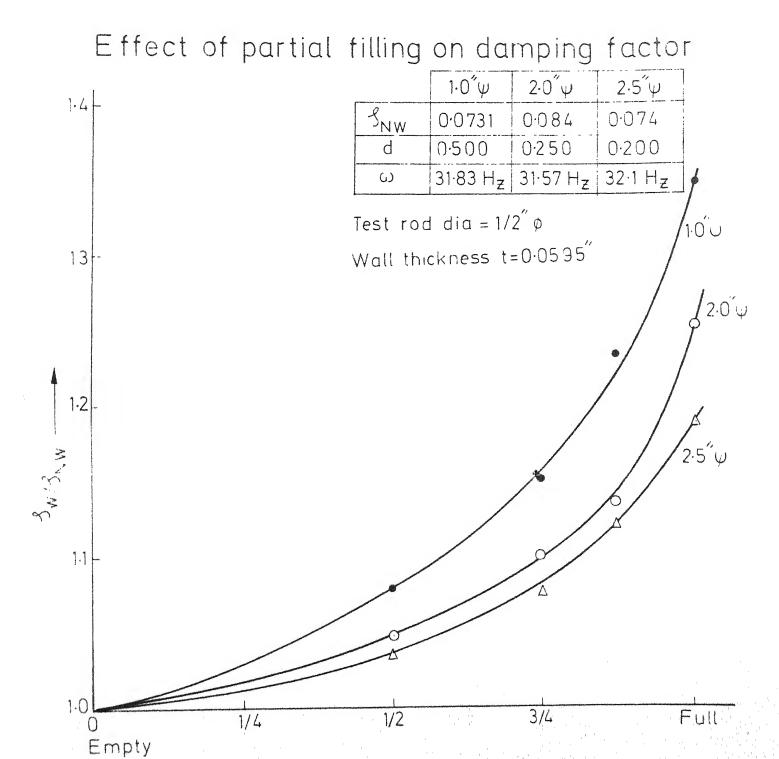
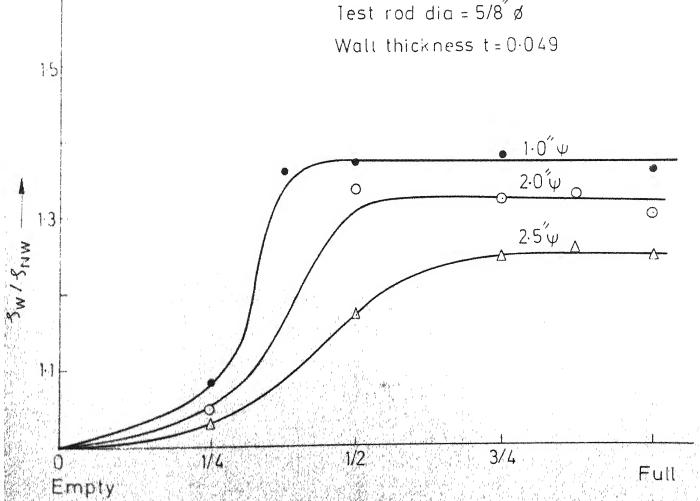


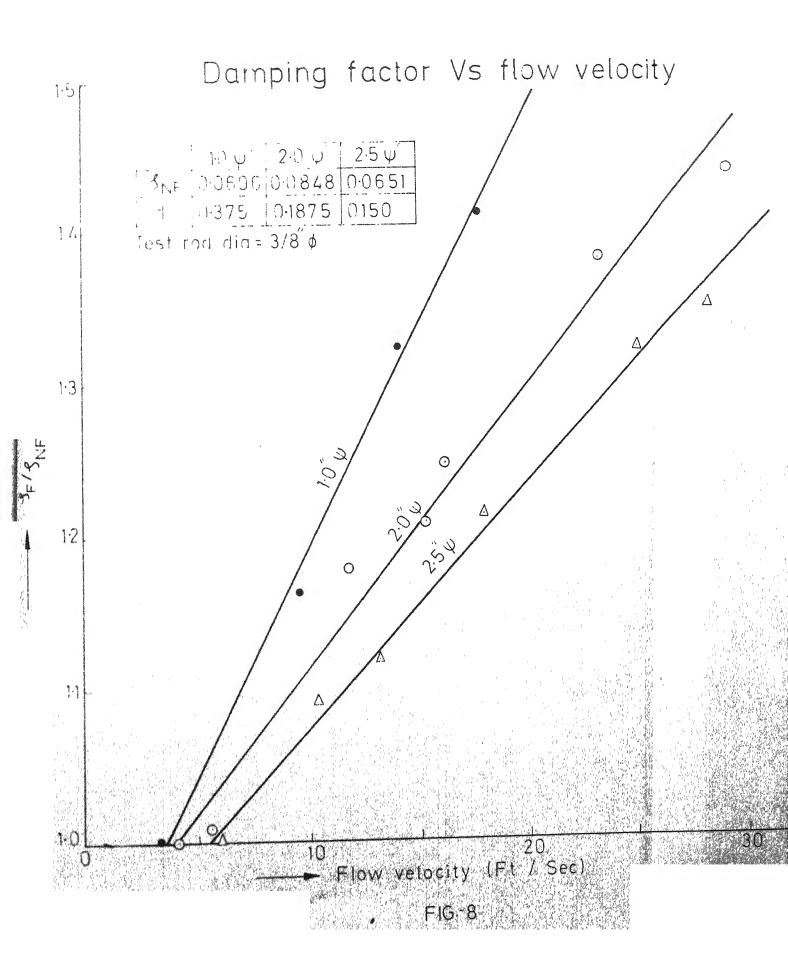
FIG.-6

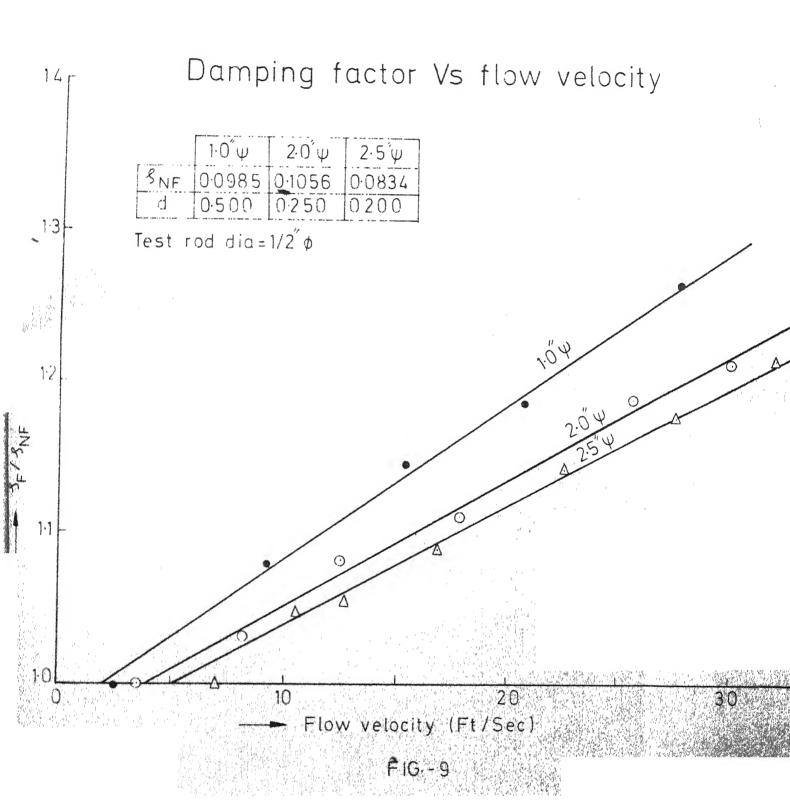
Effect of partial filling on damping factor

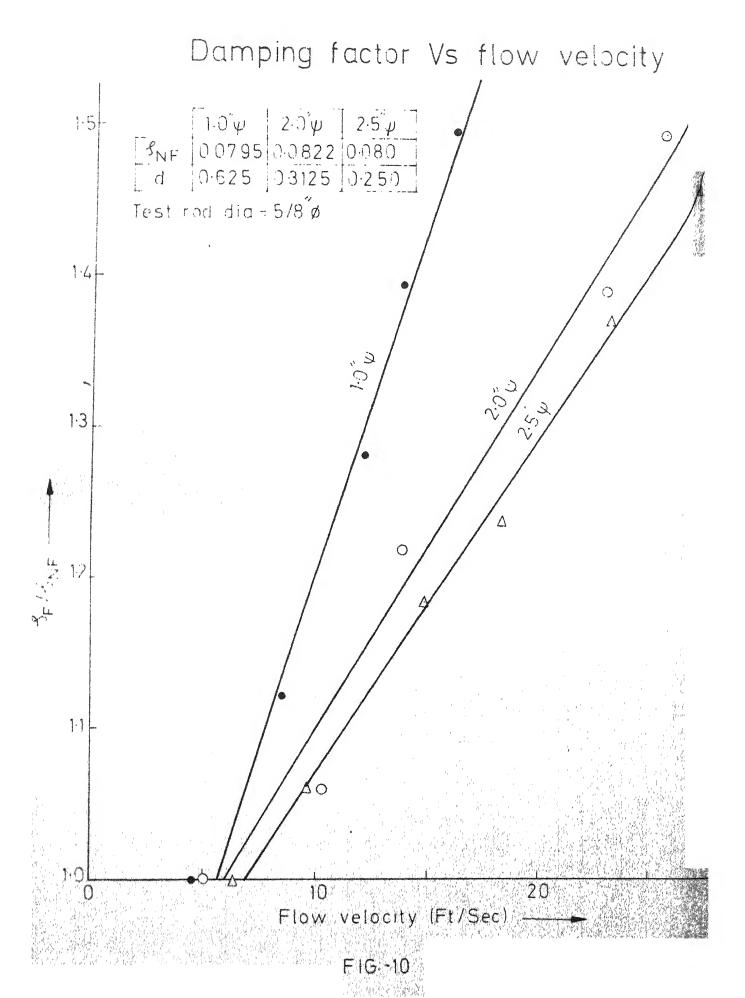
	1.0″ψ	2.0″ψ	2.5 ψ
SNW	0-058	0.0628	0064
d	0.625	0.3125	0.250
ω	37.2 Hz	36.4 H _Z	36·15 H _z

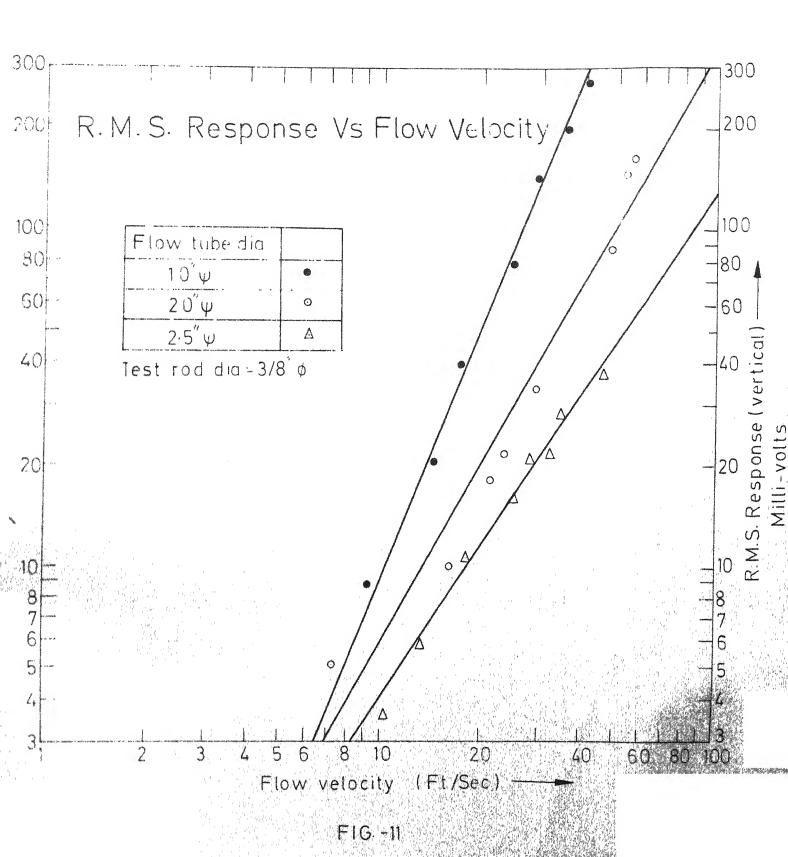
lest rod dia = 5/8 ø

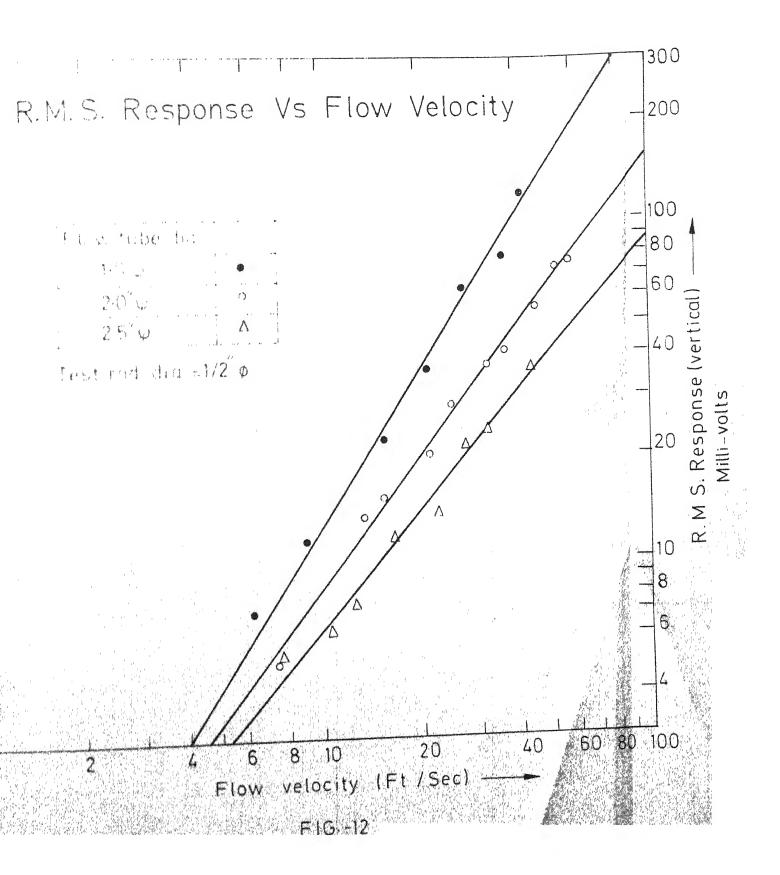


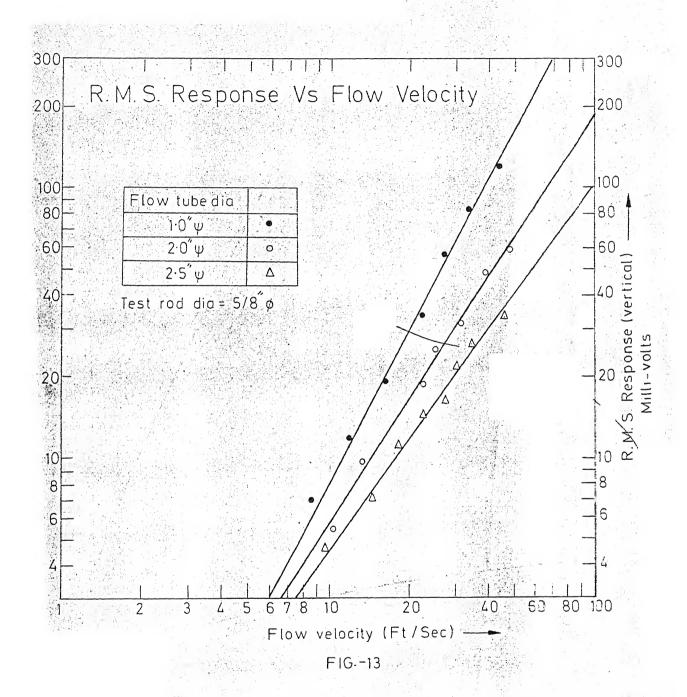


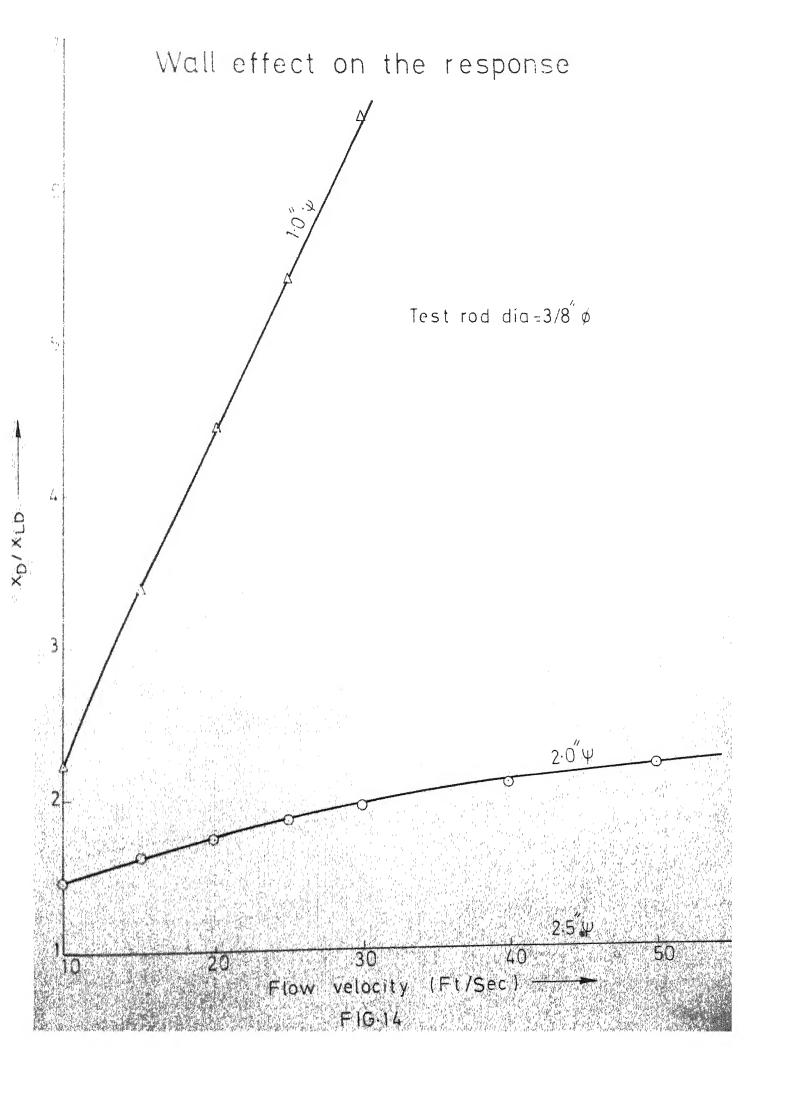


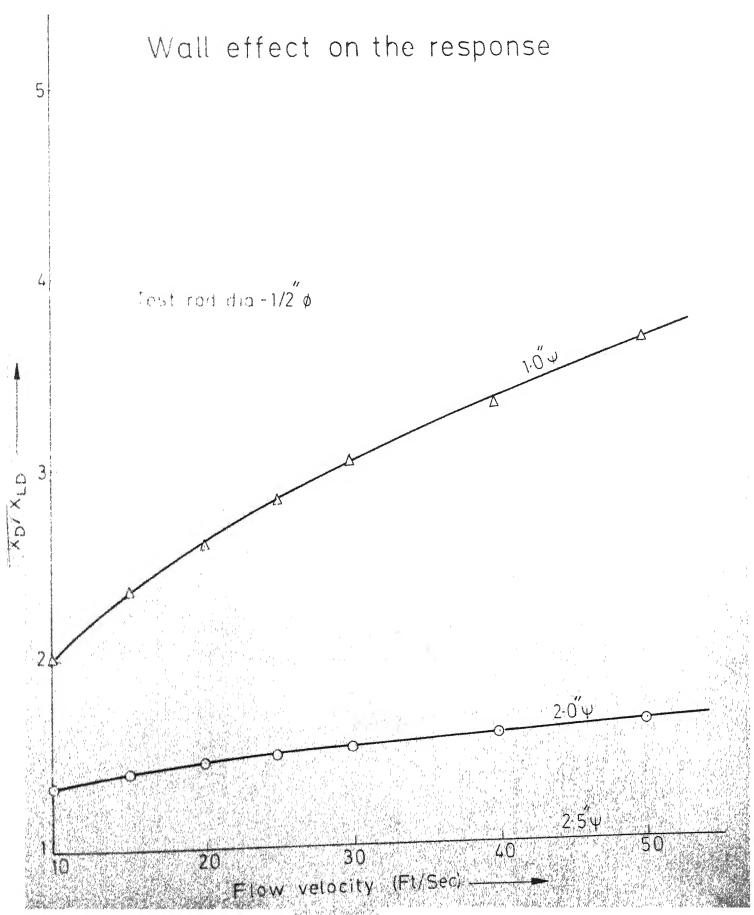












F16-15

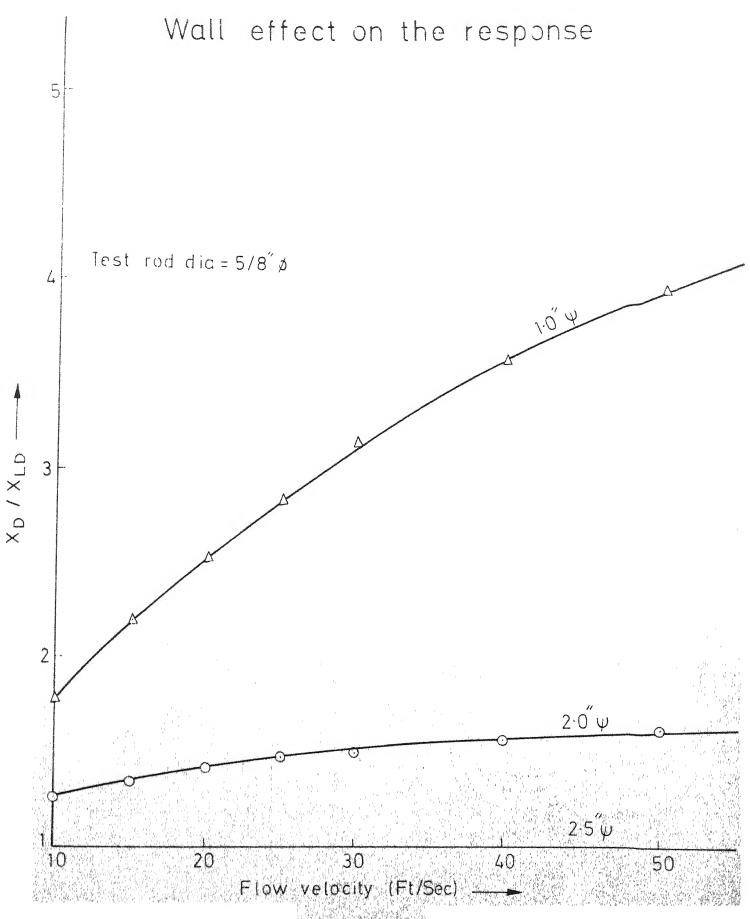
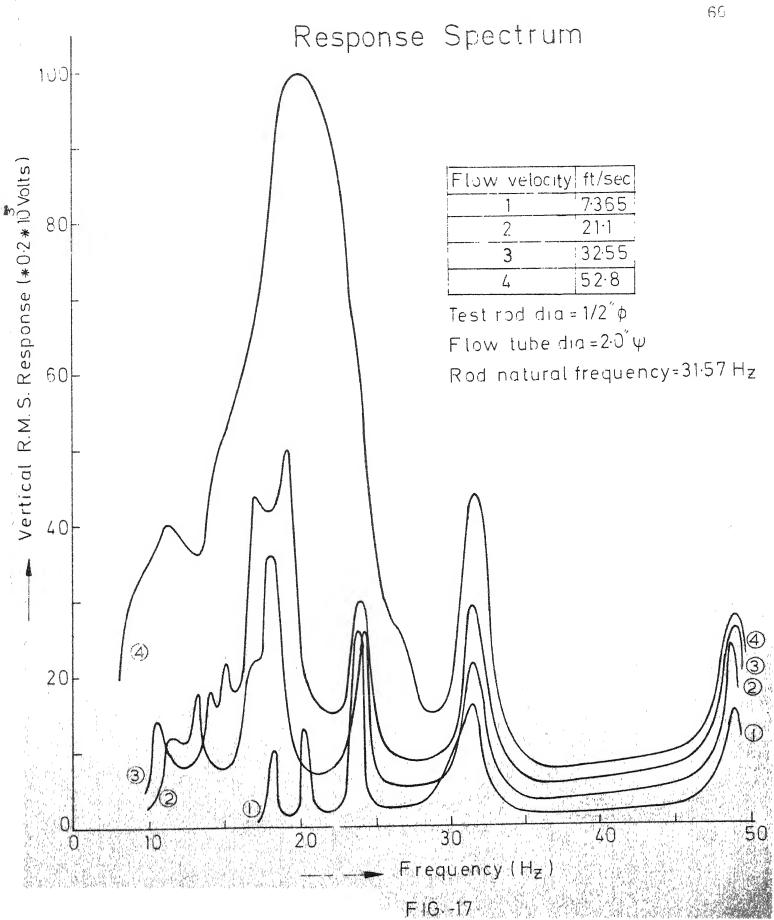
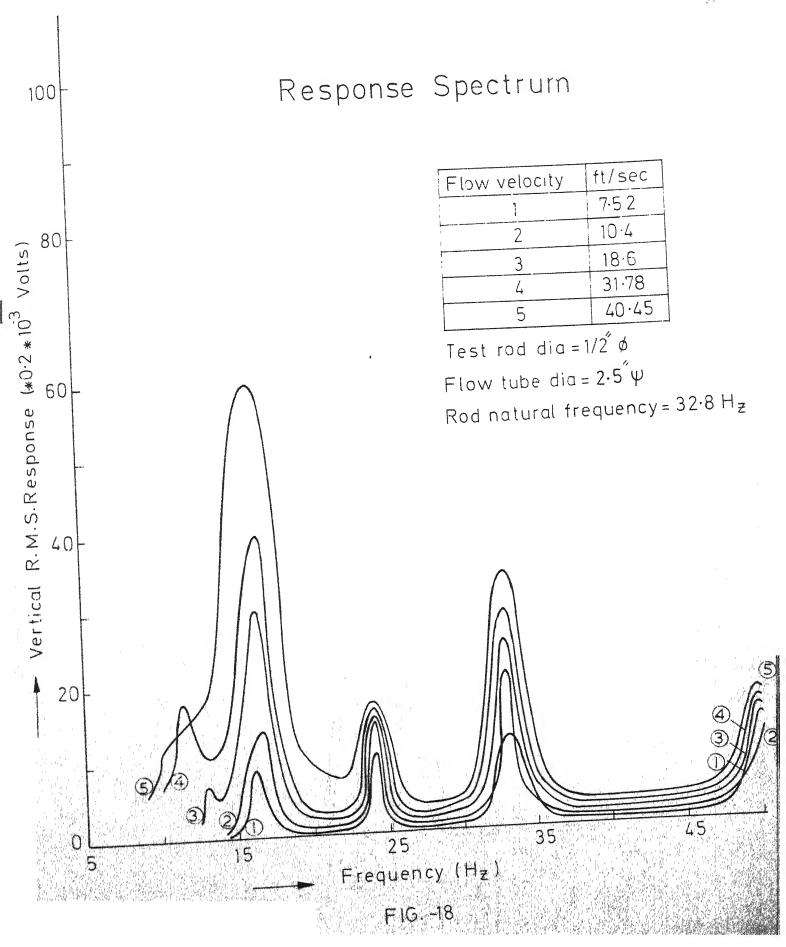
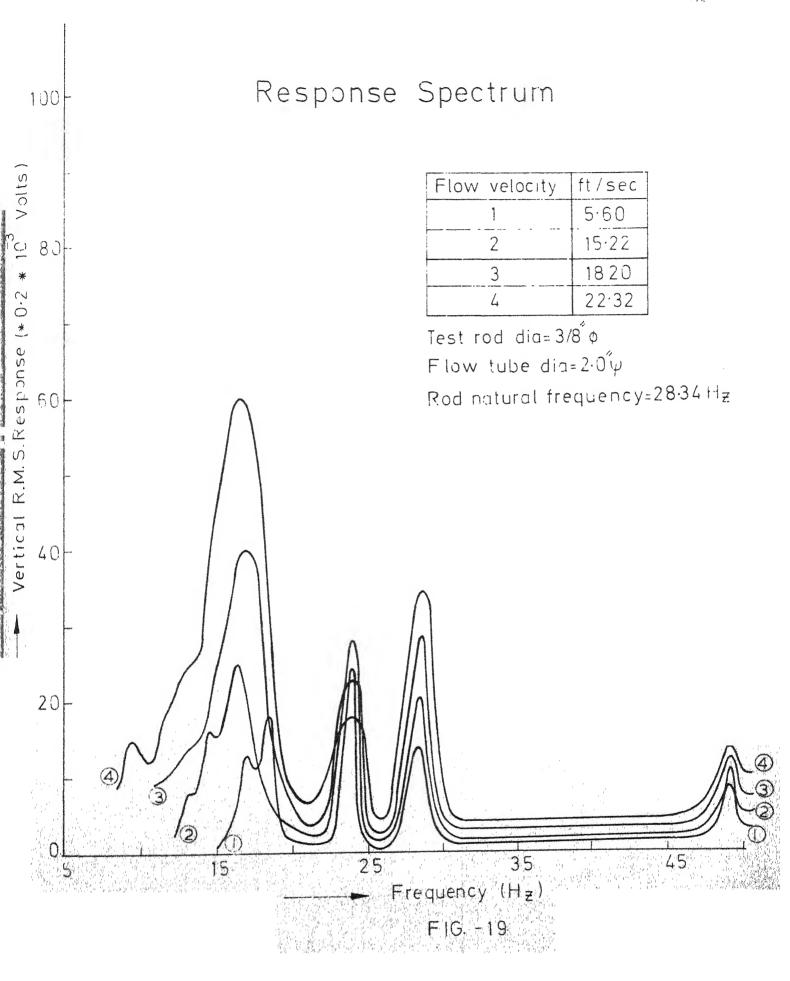


FIG-16







Response Spectrum 60 Flow velocity ft/sec 3.385 Vertical R M.S. Response ($*0.2 * 10^{-3}$ Volts) 8.13 13 47 Test rod dia=3/6° p Flow tube dia = $1.0^{\circ} \Psi$ Rod natural frequency = 29.1 Hz 40 20 45 35 25 15 Frequency (Hz)

FIG -20

CHAPTER-IV

CONCLUSIONS AND SCOPE FOR FUTURE WORK

4.1 Conclusions:

4.1.1 Effect of Water Level on the Damping:

- (1) Damping factor increases with the water level; its rate of increase, decreases for the stiffer test rod.
- (2) Rate of increase in damping and wall confinement effect will be more, if the initial damping of the test rod (with no-water) is small.
- (3) Damping contribution due to normal drag force will be more for the larger diameter test rod and will result in an increase in damping upto 1/2 water level. After that, it becomes more or less constant with the increase in water level in the flow tube.

4.1.2. Effect of Flow Velocity on the Damping:

- (1) Damping increases with the flow velocity and follow a straight line fit.
- (2) For the lower stiffness test rod, wall effect is prominent even at lower values of wall confinement parameter d.

- (3) If the damping at the initial flow velocity is large, the slope of damping line is small and wall confinement has not much effect on the increase in damping.
- (4) Natural frequency of damped oscillations remains almost constant for the moderate velocity range (0 to 30 ft/sec) and also the wall confinement has very little effect on it.
- (5) Contribution due to normal drag force, with the increase in flow velocity, helps in raising up the damping values for the larger diameter test rod.

4.1.3 Effect on the R.M.S. Response:

- (1) R.M.S. response increases with the flow velocity and follow a power function relationship. The exponent of the power function relationship ranges between 2.5 to 1.15 in the present study. Exponent values will depend on the flow noise, damping, rod end-support conditions and flexural rigidity of the test rods.
- (2) Response increases with the wall confinement for the same test rod in various flow tubes. Wall confinement has greater effect on the response level, if the exponent of power function relationship is large for the test rod in the largest diameter flow tube.
- (3) If the stiffness of two test rods are not far apart, the response level will be mainly governed by the wall confinement parameter d.

4.1.4 Effect on the Power Spectral Density:

- (1) At lower flow velocities, energy is mainly distributed around the rod natural frequency and some energy is available at lower frequencies. But at higher flow velocities, most of the energy density is shifted towards the lower frequencies. This is possibly due to large damping of the test rod (wide frequency-response function) and the turbulence decay process of the larger eddies (resulting in shifting of the energy density of the forcing function to lower frequencies)
- (2) Wall confinement tends to further increase the energy density at lower frequencies as compared to rod fundamental frequency.

4.2 Scope for Future Work:

In the present work, the wall confinement effect was studied for the test rods having larger stiffness for larger diameter. These studies resulted in high values of response for smallest stiffness test rod and thereby giving the impression of larger wall effect, though the wall confinement parameter d is small as compared to larger diameter test rods. It will be worthwhile to work on the test rods having lower stiffness with larger diameter (to know that how the disturbances help in the growth of response, as the wall confinement increases) and constant stiffness with different diameter test rods (to investigate the absolute effect of wall confinement on the response level).

Wall confinement effect can be understood in a better way, if we have the knowledge of the forcing function distribution (random pressure measurement in the turbulent boundary layer around the test rod) and the frequency-response function for the test rods (mounted in the external flow tube).

The response sensing device should be such that it should not disturb the flow at the point where it is attached to the test rod (unlike the strain gages used in the present study). One of the solution is to use miniature piezoelectric transducer flush mounted in a cavity in the test rod wall thickness.

Above stated is the work mainly connected with the vibrational aspect of the present problem. Deep understanding of the problem can only be achieved by studying also the fluid mechanics aspect. It will be interesting to know that how the flow disturbances are originated, grow-up and cause the test rod to vibrate; how the vibrating test rod affects the disturbances producing it and what is its consequence on the vibration amplitude when the wall confinement is increased. This will further require to understand the mechanism of energy transfer from flowing fluid to the solid.

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APPENDIX-1

Equation of Motion as Obtained by Chen & Wambsganss (23):

An small element &x of the rod in axial flow, with velocity V, was considered to be acted upon by axial tension T, shear force Q, bending moment M, inertia force of the rod $m_r(\frac{\partial^2 y}{\partial t^2})$ x, external force on the rod lateral surface $(q. \delta x)$, $(F_D. \delta x)$ viscous damping force, $F_L \& F_N$ drag forces per unit length in longitudinal and transverse direction, F_I lateral force per unit length on the rod and is given by the change in momentum of lateral flow about the rod (12).

$$F_{I} = m_{f} \left(\frac{\partial}{\partial t} + V \frac{\partial}{\partial x} \right)^{2} y \qquad (1)$$

where $m_{f} = added virtual mass of the rod per unit length.$

$$= \frac{\pi}{4} D^2 / C_{M} ,$$

here $C_M = coefficient$ of added mass

D = diameter of rod

 β = density of fluid.

Rod material was postulated to obey stressstrain relation of Kelvin type, that is,

stress,
$$\sigma = (\mathbb{E} + \mu \dot{e})$$
 . . . (2)

where E = Young's modulus of elasticity

/ = Internal damping coefficient

e = Strain and & = time derivative of strain.

Translatory and rotational equilibrium of the rod element leads to:

$$EI \frac{\partial^{4}y}{\partial x^{4}} + \mu I \frac{\partial^{5}y}{\partial x^{4}} + m_{f} \left(\frac{\partial^{2}t}{\partial t} + V \frac{\partial^{2}y}{\partial x}\right)^{2}y + F_{N} + F_{D}$$

$$+ F_{L} \frac{\partial^{2}y}{\partial x} - T \frac{\partial^{2}y}{\partial x^{2}} + m_{r} \frac{\partial^{2}y}{\partial t^{2}} = q \qquad (3)$$

Drag forces for rough cylinder, as discussed by Taylor (33), are taken to be -

$$F_{N} = \frac{1}{2} \rho DV^{2} (C_{D} \sin^{2} \theta + C_{f} \sin \theta)$$
 (4)

$$F_{T_{l}} = \frac{1}{2} \beta DV^{2} C_{f} \cos \theta \qquad ... \qquad (5)$$

where C_D and C_f are drag coefficients due to pressure and shear forces. Angle θ is related to the normal and axial components of flow velocity by -

$$\Theta = \sin^{-1} \left[\frac{U}{V} \right] ,$$

where the normal component of velocity, $U = (\frac{\partial y}{\partial t} + V \frac{\partial y}{\partial x})$ The longitudinal tension is the sum of the applied tension (external) and the one arising from the fluid friction. For a rod, supported at both ends, the initial tension (axial) is taken as T_0 when V=0; and as V increases, no further motion of the support is allowed. For a rod with one end free, the initial tension is zero for V=0, but as V increases, the free end is subjected to a tensile force -

$$T(1, t) = \frac{1}{2} C'_{T} m_{f} V^{2} \dots$$
 (6)

where $C_{\rm T}$ is the form drag coefficient at the free end. Substitution of eqn. (5) into the equation of translatory equilibrium and subsequent integration will give the expression for the axial tension as -

$$T (x,t) = T_{o} + \frac{1}{2} C_{T} \frac{m_{f} V^{2}}{D} \left[(1 - \frac{1}{2} \gamma) 1 - x \right] + \frac{1}{2} (1 - \gamma) C_{T} m_{f} V^{2}$$
(7)

where $\gamma=1$, if the downstream end is supported such that the displacement is zero; and $\gamma=0$ if it is unsupported or elastically supported.

The viscous damping force can be expressed as $F_D = C \frac{\partial y}{\partial t} \dots$ (8)

where C is the effective viscous damping coefficient.

On substitution of eqns. (4), (5), (7) and (8) into eqn.

(3), will result in -

$$EI \frac{\partial^{4} y}{\partial x^{4}} + \mu_{I} \frac{\partial^{5} y}{\partial t \partial x^{4}} + m_{f} V^{2} \frac{\partial^{2} y}{\partial x^{2}} - \mu_{o} \frac{\partial^{2} y}{\partial x^{2}}$$

$$- \frac{1}{2} G_{T} \frac{m_{f} V^{2}}{D} \left[(1 - \frac{1}{2} \gamma) 1 - x \right] \frac{\partial^{2} y}{\partial x^{2}}$$

$$- \frac{1}{2} (1 - \gamma) C_{T} m_{f} V^{2} \frac{\partial^{2} y}{\partial x^{2}} + 2 m_{f} V^{2} \frac{\partial^{2} y}{\partial x \partial t}$$

$$+ \frac{1}{2} (C_{N} + C_{T}) \frac{m_{f} V^{2}}{D} \frac{\partial^{2} y}{\partial x} + \frac{1}{2} C_{N} \frac{m_{f} V}{D} \frac{\partial^{2} y}{\partial t}$$

$$+ (m_{r} + m_{f}) \frac{\partial^{2} y}{\partial x^{2}} = q (x , t)$$

$$(9)$$

The Boundary conditions of the system are:

for
$$x = 0$$
, $k_1 y + EI \frac{3 \cdot 3y}{3 \cdot x^3} = 0$ and $C_{10} \frac{3y}{3x} - EI \frac{3^2y}{3x^2} = 0$

for $x = 1$, $k_2 y - EI \frac{3^3y}{3x^3} = 0$ and $C_{20} \frac{3y}{3x} + EI \frac{3^2y}{3x^2} = 0$

(10)

where C_1 , C_2 are the torsional stiffness and k_1 , k_2 are displacement stiffness at two ends of the rod.

APPENDIX-2

Part (A) Diffuser Dimensions:

Diffusers were used to connect the 6 inch pump line to the various size of test section.

Test Section diameter (inches)	Length of diffuser (inches)	(Semi Cone Angle (degrees)	Č Č
1.0	14.18	10	
1.5	12.75	10	
2.0	22.85	5	
2.5	20.0	.5	

Part (B) V-Notch Characteristics:

Flow velocity in the test section was calculated from the head measurements over the V-notch, using the following relation:-

$$Q = C_{d} \cdot (H)^{2.481}$$

where Q = Discharge in cubic feet/sec.

H = Head, over the crest, in feets.

 C_d = Coefficient of discharge.

For the V-notch used in the present study, value of the C_{d} was 2.3974.

APPENDIX-3

Typical Values of R.M.S. Response

The horizontal and vertical component of the response are given below in the Tables 1 and 2.

Table: 1

Test-rod diameter = 1/2 inch

Flow-tube diameter = 2.5 inch

System amplifier = 5000 gain

Horizontal Vertical Velocity component 0 component Sl. (ft./sec.) (Milli-volts) ≬ (Milli-volts) No. r.m.s. r.m.s. . 1. 7.52 4.8 5.3 6.4 2. 10.4 5.0 7.6 3. 12.6 6.0 9.5 11.8 16.78 4. 12.0 14.0 5. 22.45 22.0 6. 18.0 27.75 7. 31.86 22.5 24.3 36.4 8. 44.75 32.0

Table: 2

Test-rod diameter = 1/2 inch

Flow-tube diameter = 1.0 inch

System amplifier = 5000

gain

Sl.	<pre>Velocity (ft./sec.)</pre>	Horizontal component (Milli-volts) r.m.s.	<pre>Vertical component (Milli-volts) r.m.s. </pre>
1.	6.18	6.0	7.1
2•	9.2	8.5	11.6
3.	15.34	18.0	22.8
4.	20.7	32.0	, 36.5
5.	27.6	56.0	63.0
6.	36.5	61.5	78.0
7.	42.1	89.5	115.0

These tables also indicate that the wall confinement influences the horizontal signal in the similar way as the vertical signal.

ERRATA

	written as	to be read as
Page 5 , line 11	Stream	Steam
Page 7 , line 15	These	The
Page 18, line 4	16	<u>16</u>
Page 19, line 12	Circuitory	Circuitry
Page 25, line 20	hallow	hollow
Page 49, line 21	1.2 inch	1.0 inch